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1	Retrieving mean volumetric properties of multiphase flows from 2D images: a new approach combining deep learning
2	algorithms and 3D modelling
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7	Abstract
8	Measuring the morphological properties of complex multiphase systems is a crucial problem in many areas of
9	science and industry and is particularly difficult in dense environments with limited optical access. This paper
10	presents a new approach capable of extracting three-dimensional (3D) information from spherical particle
11	systems based solely on two-dimensional (2D) projections of the system. Synthetic images of the system are
12	generated using a stochastic geometrical model from a simulated 3D particle system with the same geometrical
13	features as the studied system, which is projected into 2D images labeled with the appropriate 3D information.
14	These images are then fed to a convolutional neural network (CNN) for training before being tested on synthetic
15	and experimental images. Validation results show that this technique successfully predicts the mean features of
16	the studied systems, even for dense environments with overlapping particles, with high computational efficiency.
17	
18	Keywords: Convolutional neural network, Particle systems, Stochastic Geometry, 3D Modelling.
19	
20	1 Introduction
20	

Complex particulate systems are an important and widely studied feature of many industrial processes (Black et
al., 1996; Gianinoni et al., 2003; Honkanen et al., 2010) and research projects (Clift, R., Grace, J. R., & Weber,
2005; Poelma, 2020). These systems are generally defined as mixtures of two or more substances where one is
suspended in another, either a gas in a liquid (bubble flow) (Juliá et al., 2005; Karn et al., 2015; Lau et al., 2013),
a solid in a liquid (Kavanaugh et al., 1980; Yu et al., 2009), or one immiscible liquid in another (emulsions) (Huang
et al., 2001; Maaß et al., 2011).

27 Measuring the properties of these systems is essential to optimize and improve the performance of many processes 28 involving multiphase flows (Emmerich et al., 2019; Panckow et al., 2017). In chemical engineering for example, 29 knowledge of particle spatial distributions is crucial to calculate mass and heat transfer rates and the reaction 30 kinetics governing the efficiency of the process. In this context, the most important properties are the particle size 31 distribution (PSD) of the dispersed field, the mean characteristic diameters (d_{43} , d_{32} , d_{10} , ...), the volume fraction 32 of the dispersed phase (ϕ), and morphological information such as the shape and irregularity of the particles.

33 Extracting this information is a non-trivial problem and many approaches have been investigated in the literature. 34 Laser-based methods involve the analysis of scattered light from a laser beam passing through the system. One of 35 the most popular in-line method is focused beam reflectance measurement (FBRM) (Heath et al., 2002; Ruf et al., 36 2000), in which a highly focused rotating laser beam is passed at a fixed speed over the suspended particles and 37 the duration of the backscattered light is measured. Although this only provides information on the chord length distribution of the particles, and post processing is required to retrieve the PSD, this technique has successfully 38 39 been used to measure droplet size distributions in water oil emulsions (Boxall et al., 2010) and micro-bubble size 40 distributions in air flotation processes (Couto et al., 2009), and has proven particularly valuable for the 41 characterization of crystal-like particles (Acevedo et al., 2021; Heinrich and Ulrich, 2012; Pandalaneni and 42 Amamcharla, 2016; Pandit et al., 2019). The second widely used laser-based technique in this context is digital in-43 line holography (DIH), otherwise known as lens-free imaging (Darakis et al., 2010; Lamadie et al., 2012), in which 44 particle characteristics are estimated from the laser diffraction patterns of the system. This approach has been used 45 to study the position and size of particles in pipe flows (Sentis et al., 2017), microscopic setups (Sheng et al., 2006), and sprays (Yang and Kang, 2011), combined with a machine learning algorithm to study non-uniformly 46 47 shaped particles (Shao et al., 2020). However, while DIH is very efficient for 3D positioning and PSD 48 measurements, it can only be used in optical dilute media.

49 Imaging-based methods generally provide richer information than other available techniques and therefore more 50 widely used. Image analysis algorithms combined with direct imaging can be used to measure properties other 51 than the PSD (Maaß et al., 2011), such as the volume fraction of the dispersed phase (Karn et al., 2015), 52 morphological information on irregular particles (Suh et al., 2021), and the presence of clusters (Zhang et al., 2012). The image processing algorithms used have included deterministic methods such as the Hough transform 53 54 (Yu et al., 2009), watershed segmentation (Chen et al., 2004), more sophisticated approaches to deal with 55 overlapping particles (de Langlard et al., 2018a; Zafari et al., 2020; Zhang et al., 2012; Zou et al., 2021), and 56 recently, machine learning methods (Cui et al., 2022; Haas et al., 2020; Kim and Park, 2021; Li et al., 2021). Most of these techniques extract 2D information from the detection of individual particles and cannot measure 3D properties. Although 3D information can be retrieved from multiple viewpoints, this requires a complex imaging setup with multiple optical access points (Wang et al., 2022; Xue et al., 2014), which is not always desirable or possible.

This article presents a new approach in which a deep learning algorithm is combined with stochastic geometrical models to extract 3D properties from 2D projected images of the system. The 3D stochastic geometrical model is used to generate a huge set of synthetic 2D images labeled with the 3D geometrical properties of the particle field and with the same geometrical properties as experimental images. This model can be used to reproduce any 3D particle field and generate 2D projections suitable to train a convolutional neural network (CNN) (Dia et al., 2022). To the best of our knowledge, this type of approach has only previously been used by Fend et al. (Fend et al., 2021) to reconstruct highly porous 3D structures from focused ion beam scanning electron microscopy data.

The paper is divided into six sections. The following methods section provides a brief reminder of the principles of stochastic geometrical modelling and a description of the chosen machine learning algorithm. The third section describes the parameters chosen to measure the performance of the network and the results of tests on representative simulated flows. The fourth section presents experimental results for 3D dispersed phase volume fractions and PSDs retrieved from highly concentrated particle systems. The results obtained for experimental and liquid-liquid system images are presented in section 5 and the final section is an overall conclusion with perspectives.

75 **Computational methods**

As mentioned above, the proposed approach (Figure 1) involves two computational tools, a stochastic geometrical
model and a CNN.

For the synthetic images to be representative, the simulated particle system should have the same geometrical properties as the studied system, i.e. be statistically representative of the observed particle field. The 3D stochastic model used to generate the images should consist of a hard-core model of the particles, in this case spherical, that accounts for the statistical properties of the dispersed phase by eliminating all particle-particle and particle-wall interactions. Matérn type II point processes (Matérn, 2014) were therefore chosen to generate the synthetic images. Matérn type II point processes are derived from an underlying homogeneous Poisson point process in \mathbb{R}^d with

84 intensity λ . A point process is simulated and a thinning rule is applied based on a hard-core distance (r > 0)

85 between the generated points, which removes the last arriving point from any pair of points less than 2r apart. This process can be considered a marked point process where the first mark r (constant and positive) represents 86 87 the hard-core radius and the second mark is the time of arrival of the points modeled using a uniform random 88 variable. 89 The generation process for the synthetic images thus involves the following steps: 1. Choosing a 3D domain $W = \mathbb{R}^2 \times [0, l]$, where l > 0 is the length of the projection direction, consistent 90 91 with the dimensions of the actual measurement volume. 92 2. Randomly positioning the particles inside W according to a Poisson point process with intensity λ . Each 93 point is marked with a radius r_i , and a time of birth $t_i, \forall i \in \{1, ..., n_p\}$. 94 3. Applying two thinning rules: i) Matérn's thinning rule, i.e. eliminating the last arriving point

95 $(\max(t_1, t_2))$ of pairs (x_1, r_1, t_1) (x_2, r_2, t_2) with $||x_1 - x_2||_2 \le r_1 + r_2$, and *ii*) de Langlard et al.'s thinning

96 rule (de Langlard et al., 2018b) to eliminate particle–boundary interactions by removing points with a

97 probability $1 - \exp(-U(x, r, t))$, where $U: W \times \mathbb{R}^+ \times [0, 1] \to \mathbb{R}^+ \cup \{+\infty\}$ is an interaction function.

98 4. Projecting the generated 3D particle field orthogonally onto a 2D grid to create a single synthetic image, as99 shown in Figure 1.

100 The model generates continuous-valued particles that are later discretized when the synthetic images are generated;
101 the resolution of the images can thus be adjusted as required. The whole process is repeated *N* times to assemble
102 a training database, and the geometrical properties of each 3D model (size distribution, spatial distribution, etc.)
103 are all stored along with the 2D images (cf. Figure 1, second step). Further details about the stochastic 3D model
104 used here can be found in de Langlard et al. (de Langlard et al., 2018b).

105 This model can also be used when the distances between points are non-deterministic, and is therefore applicable 106 to real systems. As demonstrated by Stoyan and Stoyan (Stoyan and Stoyan, 1985), the thinning rules are 107 generalizable to cases in which the *r* marks follow a given probability law, and generalized expressions can always 108 be obtained for distribution parameters such as the retention function and the intensity after thinning. The model 109 can also be extended by changing the shape of the particles.

1. Acquisitions and preprocessing





Figure 1. Schematic outline of the main steps of the workflow. 1- Acquisition and preprocessing of experimental images. 2 Generation of a simulated dataset and training of the neural network. 3- Prediction of 3D properties from experimental data.

Note that the training images must be of the same type as the experimental ones, and since a stochastic model is used, only binary images can be used to train the CNN. This means that the experimental data have to be acquired with a backlight setup based on a telecentric lens (cf. section 4) to limit blurring and perspective effects. The images also have to be binarized and denoised before being processed by the CNN (see Figure 1).

117 The CNN is the second pillar of the proposed approach. Convolutional neural networks are a class of deep learning 118 algorithms designed to process data with a grid-like topology (*e.g.* images, grid cells, financial series...). These 119 algorithms are classed as supervised in the sense that they require both data and labels to be trained to perform a 120 given task. They can be used to solve classification and regression problems but here, only regression CNNs were 121 considered, because the problem involves continuous values.

- **122** A CNN broadly consists of four layers (Figure 2):
- An input layer that processes the data provided to the network.

- 6
- A feature extraction or convolutional layer, which combines convolutional operations, a non-linear
 activation function (*e.g.* ReLU, sigmoid), and a pooling operator (*e.g.* max pooling, average pooling) to
 automatically extract relevant features from the images.
- A fully connected (FC) layer at the end of the network consisting of multiple interconnected neurons
 whose weights are adjusted to teach the network to identify important features.
- An output layer, which presents the predictions.
- 130 Unfortunately, there is no generic way to determine the optimal network shape (*e.g.* the number of convolutional
- 131 layers, number of neurons, number of layers in the FC layer, the learning rate) in advance. This problem is
- 132 typically solved using prior experience of CNNs for initial guesses and empirical trials to choose the best
- 133 architecture. Here, two separate networks were used, one to determine the volume fraction and one to measure
- the PSD.



138

136Figure 2. Visualizations of the architectures of the two CNNs used in the study: one to predict the volume fraction $(CNN^{\phi},$ 137upper row) used and the other to determine the PSD $(CNN^{PSD},$ lower row). The plots were generated with PlotNeuralNet

(Iqbal, 2018).

The first network, CNN^{ϕ} (Figure 2, top panel) is used to predict the volume fraction (ϕ) of the dispersed 139 1. 140 phase and was trained on volume-fraction labeled images. The input layer consists of three identical 256 × 256 images. The feature extraction layer consists of six convolutional layers respectively containing 16, 141 142 32, 64, 128, and 256 filters, with a kernel size of 3×3 . These layers also contain an activation function, 143 ReLU, are followed by a 2×2 max pooling layer, producing flattened feature maps that are fed to the FC layer. The FC layer contains three sub-layers with 512, 128 and 64 neurons, respectively, and a dropout of 144 145 50% is applied after each layer to reduce overfitting. Finally, the output layer consist of a single neuron representing the predicted value, $\hat{\phi}$. 146

147 2. The second network, CNN^{PSD} (Figure 2, bottom panel), used to predict the 3D PSD, has the same

architecture except for the output layer which is a vector of size *n* (note that the output size can be adjusted

depending on the studied system). Since the estimated parameter is a distribution, the sum of the elements in

150 the vector is normalized to 1 using the softmax function.

151 **2. Network training strategy**

152 After setting the number of layers and the parameters of the CNN, optimal hyper-parameters were determined by gradient-based optimization. The dataset (N = 22400) used to train the network was split into three categories; 153 154 an initial training set ($N_t = 16\ 000$), a validation set ($N_v = 2\ 000$) used to monitor the overall performance of the network at each iteration, and a final testing set ($N_e = 4400$) used to evaluate the performance of the 155 network. All the generated images corresponded to a measurement volume of $25.79 \times 25.79 \times 53$ mm³ in 156 157 keeping with the experimental setup. The parameters of the stochastic model were varied uniformly to account 158 for the variety of possible experimental images in real life applications. The intensity of the particles, λ , was adjusted between 100 to 250, while the PSD was modeled as a truncated lognormal function, \mathcal{L}_{tr} (μ , σ^2 , a, b), 159 160 with μ varied randomly from -7.5 to 0.06, σ from 0 to 2.5, and a and b equal to 0.005 and 0.08 respectively 161 (check Table 1).

As mentioned above, CNN training datasets have to be labeled with the target information. Here, all the numerical images were generated over a 5.5 h period and each image, *i*, was labeled with the corresponding 3D volume fraction $\phi_i = \frac{v_i}{v_i}$ (with v_i is the volume of the particles, and V_i is total volume) and number weighted PSD (PSD_i) across *n* classes.

166

Table 1. Details on the synthetic dataset.

Dataset	
16000	
2000	
4400	
[0.5 12.2]	
256 × 256 × 3	
	Dataset 16000 2000 4400 [0.5 12.2] 256 × 256 × 3

167

169 The hyper-parameters were then optimized using mean absolute error (MAE) as the loss function. The MAE

170 measures the average absolute difference between true $Y = \{y_1, \dots, y_{N_t}\}$ and predicted values $\hat{Y} = \{\hat{y}_1, \dots, \hat{y}_{N_t}\}$:

171
$$MAE(Y, \hat{Y}) = \frac{1}{N_t} \sum_{i}^{N_t} |y_i - \hat{y}_i|,$$

Finally, a series of tests were conducted to identify the best number of epochs, the optimal learning rate, and the
size of the mini-batches. Slight differences in computational times were observed, so these parameters were
optimized as follows:

175 1. Number of epochs: the CNN was trained using mini-batches due to limited computational resources, with

176 one epoch corresponding to one training run over all the mini-batches. The network training convergence

177 was reached after 100 epochs, therefore the number of epochs was then fixed at 150 with a callback to stop

the training run if the MAE increased for 10 straight epochs, to avoid overfitting.

Learning rate: The learning rate was adjusted from a starting value of 10⁻⁴, with a decay rate of 20% every
 50 epochs, down to a final learning rate of 4 × 10⁻⁶,

181 3. Mini-batch size: the mini-batch size was set to 32.

182 The CNNs were trained using a supercomputer cluster equipped with GPUs each with 640 tensor cores and a

183 memory of 32 GB. The training processes took about 30 min for each CNN.

- 184 Since the accuracy of the predictions cannot be determined for regression CNNs, the performance of the
- 185 networks was determined using the MAE and the relative error (RE) for the volume fraction predictions, and
- using the MAE only for the PSD. Since the volume fraction is a scalar quantity, MAE calculations are
- 187 straightforward, but for the PSD, which is a vector, the overall MAE, \overline{MAE} , was calculated by comparing the
- 188 average values of PSD_i and $\widehat{PSD_i}$.

189 3. Numerical validation

- 9
- 190 The performance of CNN^{ϕ} was evaluated using the testing dataset, with hold-up values ranging from 0.5 to
- 191 12.2%, and the PSD varied between images.
- 192 The results of the predictions compared to the ground truth are shown in Figure 3. The MAE of 0.21 ± 0.26
- 193 highlights the overall accuracy of the network's predictions.



Figure 3. Scatter plot of the predicted 3D volume fraction φ̂ as a function of the ground truth φ. The insets show examples of
the images from the testing dataset, with φ_(a) = 1.3%, φ_(b) = 5%, and φ_(c) = 11.3%, respectively. The points are color
coded according to the relative error of the predictions, RE_i = φ_i-φ_i, green if |RE_i| ≤ 0.05, orange if 0.05 < |RE_i| ≤ 0.1,
and red if |RE_i| > 0.1. The dashed lines border the predictions with less than 10% RE and the first bisector is represented by
a solid line. The shaded green area represents the 95% confidence interval of the predictions.

The density of particles in the images naturally increases with the particle volume fraction, making the images more complex to process (compare insets (c) and (a) in Figure 3, for instance). Significant patterns that the network uses to refine its predictions may also disappear. Nevertheless, while the relative errors of the predictions increase with the particle volume fraction, they are all of the correct order of magnitude, confirming the robustness of the approach. Figure 3 shows for instance that below $\phi = 10\%$ (93% occupied area in the images), more than 77% of the network predictions are within the 95% confidence interval. 206 To validate the PSD estimates, 50 further sets of 100 images were generated using a truncated lognormal function with a fixed particle intensity of $\lambda = 200$. The variance s^2 and the mean diameter \bar{d} were chosen to 207 208 cover a wide range of narrow and wide unimodal distributions with mean diameters ranging from 500 µm to 209 5 mm. Predictions were made for each image using the trained CNNPSD and binned in 13 size classes. The 210 predictions were then averaged (\overline{PSD}) and compared to the ground truth distribution (PSD) obtained from the 211 corresponding function. 212 The \overline{MAE} values obtained (check Table 2) were lower than 0.03 in 60% of the considered cases, corresponding 213 to good agreement between the predicted and ground truth distributions (see Figure 4). For the broadest distributions (variance $\geq 10^{-3}$), corresponding to a maximum occupied surface fraction of about 80%, 82% of 214 215 the \overline{MAE} values are less than 0.03. As expected on the other hand, the network performs less well for very

- narrow distributions (variances range from 10^{-5} to 10^{-4}), as there were no such distributions in the training 216
- 217 dataset. The results for four typical cases are presented in Figure 4.
- 219

218 Table 2. Overall MAE for PSD predictions from all 50 simulations. Results for the four cases presented in Figure 4 are highlighted in bold.

Variance	s ²				
(mr	n) 10 ⁻⁵	10-4	10-3	10-2	10-1
Mean \bar{d}	10	10	10	10	10
(mm)					
0.5	0.059	0.021	0.009	0.009	0.009
1.0	0.085	0.027	0.009	0.009	0.009
1.5	0.071	0.027	0.012	0.010	0.010
2.0	0.096	0.043	0.014	0.011	0.011
2.5	0.114	0.068	0.018	0.011	0.011
3.0	0.115	0.080	0.024	0.013	0.011
3.5	0.123	0.083	0.034	0.014	0.012
4.0	0.115	0.077	0.041	0.015	0.012
4.5	0.114	0.088	0.047	0.018	0.013
5.0	0.131	0.108	0.058	0.022	0.013



221

Figure 4. Comparison of CNN^{PSD} predictions with the ground truth for a representative selection of simulated mean particle diameters and particle size variances. The inset in each panel is a typical image of the corresponding particle field. Error bars represent the standard deviations of the predictions.

225 4. Experimental validation

After validating the approach on simulated data, experimental validations were performed, first using calibrated spherical poly(methyl methacrylate) (PMMA) beads (diameter precision $\pm 1\%$) submerged in brine stirred in a rectangular mockup unit and second, using existing data from an emulsion in a cylindrical stirred tank.

229 4.1.Validation with calibrated beads

230 Particle suspensions are particularly convenient for validation as they reproduce the main features of two-phase

- flows but without the coalescence and breakage events that frequently occur with bubbles and droplets. With
- solid particles moreover, the properties of the dispersed phase are clearly defined. Here, calibrated spherical
- 233 PMMA particles were stirred in 1 L of brine (330 g NaCl dissolved in 1 L of water) using a four blade propeller.

Brine was chosen so that the density of the continuous phase could be adjusted to that of the dispersed phase to

avoid creaming effects and guarantee a homogeneous dispersion of the particles in the flow.

- A lognormal distribution of particle sizes cannot be reproduced experimentally using calibrated beads because of
- time and cost constraints (a 1.48% volume fraction corresponds to about 10000 particles for instance).

238 Therefore, only four sizes of PMMA particles were used (0.39, 0.5, 0.79 and 1.59 mm in diameter), with four

volume fractions and a fixed PSD. The number of particles used in each experiment was determined by weighing

240 them with a benchtop precision scale (\pm 0.2 mg), to evaluate uncertainties on the measured hold-up values. The

- various experimental configurations are listed in Table 3.
- Table 3. Details of the volume fractions and size distributions used for validation experiments with calibrated bead
- 243

Experiments	3D hold-up (%)	Size Distribution (PSD)			
Experiments	ϕ (± scale error)	0.39 mm	0.5 mm	0.79 mm	1.59 mm
1	1.48 ± 0.002	50%	30%	15%	5%
2	2.91 ± 0.003	50%	30%	15%	5%
3	4.31 ± 0.005	50%	30%	15%	5%
4	5.67 ± 0.006	50%	30%	15%	5%

suspensions.

244

245 The mixture was placed in a rectangular tank with external dimensions of $25 \times 15 \times 6.7$ cm³ and a capacity of

246 1.1 L, fitted with two optical quality windows. The mixture was agitated using a propeller at 2000 rpm to

247 produce a homogenous flow (Figure 5).



248

Figure 5. (a) Side-view sketch of the experiment setup with from left to right, the high speed camera, bi-telecentric lens, tank, and collimated green light-source. (b) Front view of the tank. (c) Computer with a typical example of a captured image on the screen ($\phi = 2.91\%$).

Two-dimensional projections of the 3D particle field in the tank were acquired using a backlight setup to avoid
perspective effects. A Photron FASTCAM Mini UX100 CMOS camera was attached to an Opto engineering
TC16M096 bi-telecentric lens placed 26.2 cm from the center of the tank. The tank and sensor were uniformly
illuminated with an Opto engineering collimated green light source (peak wavelength, 525 nm; beam diameter,
120 mm).



258 29.9 mm² (scale factor, 2.6×10^{-2} mm/px). The images were binarized, cropped to 980×980 squared pixels

- 259 $(25.8 \times 25.8 \text{ mm}^2)$, and resized to 256×256 squared pixels to facilitate coding. The tank was 5.3 cm wide
- along the optical axis and covered the entire telecentric field of view, such that the total measurement volume
- was 35.3 cm^3 . The images were captured at 50 frames per second, with an exposure time of 1/81920 s,
- avoiding any motion blur. A set of 4 365 images was acquired for each experiment (Figure 1, first step).

263 The images underwent three steps of preprocessing, starting with a simple binarization using an arbitrary gray 264 level threshold (TH = 20). Morphological area opening using a disk shaped element of 5 pixels was then 265 applied to denoise the images, removing features such as micro bubbles and dust particles, which appear as small 266 white dots in the background. The particle contours were then smoothed using an opening operation (erosion 267 followed by dilatation) with the same structuring element. These morphological operations produce images that 268 are similar to synthetic images (cf. Figure 6). However, small differences between binarized and simulated 269 particles can be observed on pixel-by-pixel inspection. The variance between the images can be quantified using 270 the intersection over union (IoU) ratio between the set of ground truth (synthetic) particles and the set of 271 binarized particles:

272
$$IoU = \frac{|A \cap \hat{A}|}{|A \cup \hat{A}|}$$

Where *A* and \hat{A} stand for the simulated and binarized particles respectively, \cap and \cup are respectively the intersection and union operators, and the vertical bars denote cardinality.

 0.39 mm
 0.50 mm
 0.79 mm
 1.59 mm

 Raw data
 Image: Constraint of the second second

275

Figure 6. Comparison between binarized raw data (\hat{A}) and simulated particles (A). Non-intersecting regions are colored in green for pixels in \hat{A} but not in A and in magenta for pixels in A but not in \hat{A} .

The IoU values obtained for all unconnected particles in the dataset (Table 4) show that the discrepancy between

the shapes of the modeled particles and real particles is greater for smaller particles. Moreover, Figure 6 shows

that while the model underestimates the size of small particles, larger particles are enlarged in the binarized

281 images. Regardless therefore of the training and the architecture of the machine learning network used,

discrepancies between the predictions and the ground truth can be expected, especially for the smallest particles.

- 283 Note that applying the same morphological operations to the synthetic images did not improve the agreement
- between the synthetic and binarized images.
- 285

Table 4. Intersection over union area for simulated and binarized particles of different sizes.

Particle radius size (mm)	IoU
0.39 mm	0.92
0.50 mm	0.95
0.79 mm	0.97
1.59 mm	0.97

286

287 A new training dataset was generated to retrain CNN^{ϕ} considering only four particles sizes to mimic the

288 experimental conditions. The distributions of $\hat{\phi}$ values obtained are shown in Figure 7 in separate boxplots for

each volume fraction (hold-up).

290





Figure 7. Boxplot representations of the distribution of predicted hold-up values obtained for each experimental particle volume fraction. Examples of the experimental images are shown below the corresponding hold-up value. Each boxplot consists of a red box extending from the first (Q_1) to the third (Q_3) quartile of the distribution and whiskers extending to $1.5 \times [Q_3 - Q_1]$. Outlier values beyond this range are shown as scattered points. The first bisector (solid line) and 95% confidence interval (shaded area) are shown to visualize the differences between the average prediction (numerical value on the left of each box) to the corresponding experimental hold-up ϕ .

298 These results indicate that the chosen approach was effective for all experimental images. As observed in the

numerical experiments, the predictions tend to become less reliable when the particle density in the image

300 increases. Whereas more than 25% of the predictions fell within the confidence interval at = 1.48%, this

proportion decreased to 12%, 6% and finally 1% at volume fractions of 2.91%, 4.31% and 5.67%, respectively.

302 Moreover, Table 5 shows that while the relative error was within the confidence interval ($RE \le 0.05$) for the

first two experiments, the errors were an order of magnitude larger at higher volume fractions. Figure 8

304 compares the predicted PSDs with the ground truth for the same four experiments.

305

306Table 5. Comparison between the average predicted value of the volume fraction ($\hat{\phi}$) and the actual value (ϕ); std, standard307deviation; RE, relative error; MAE, mean absolute error.

	$\phi(0/)(1 \text{ and})$	Average			
Experiments	φ (%)(± scale	surface	$\hat{\phi}$ (%) (± std)	RE	$MAE \ (\pm \text{ std})$
	error)	fraction (s/S)			
1	1.48 ± 0.002	45.6%	1.49 ± 0.19	0.008	0.15 ± 0.12
2	2.91 ± 0.003	71.8%	2.95 ± 0.31	0.017	0.25 ± 0.19
3	4.31 ± 0.005	85.4%	4.76 ± 0.44	0.105	0.52 ± 0.37
4	5.67 ± 0.006	92.4%	6.51 ± 0.39	0.148	0.84 ± 0.38



309

310 Figure 8. Comparison of predicted particle size distributions (\widehat{PSD}) with the ground truth (PSD) for the four considered 311 particle volume fractions ϕ .



313 only a slight underestimation of the proportion of smallest particles, probably because of the imperfections of the

binarization procedure. At higher volume fractions, the PSD was not accurately recovered, the proportion of

- 315 larger particles being overestimated, presumably because of particle overlap. The poorer agreement with the
- actual PSDs at higher particle volume fractions is reflected by an increase in the corresponding \overline{MAE} s (Table 6).
- 317

Table 6. \overline{MAE} for the four experimental PSD prediction.

Experiment	$\overline{\text{MAE}}$ (±std)
1	0.018 ± 0.006
2	0.019 ± 0.006
3	0.024 ± 0.068
4	0.031 ± 0.008

318

319 **4.2.Validation for a liquid-liquid system**

320 The approach was further validated using previously acquired data (Amokrane et al., 2014). Images of water 321 droplets in a emulsion with tetra-propylene hydrogen (TPH, density $\rho = 760 \text{ kg/m}^3$; viscosity $\eta = 1.26 \times$ 322 10^{-3} Pa · s; surface tension $\gamma = 43 \times 10^{-3}$ N/m) were obtained using an in-situ video camera placed inside a 323 cylindrical stirred tank reactor (internal volume, 1 L; internal height, H; Figure 9 (b)). Emulsions with hold-ups 324 of up to 5% were stirred at 500 rpm with a three-flat-blade propeller (diameter, 60 mm) placed H/3 above the 325 base of the vessel. Images were acquired after 60 min of stirring to ensure that a steady state had been reached. 326 The probe consisted of a CCD camera with LED back lighting (Figure 9 (a)). Images were obtained of the 327 emulsion flowing through a 1500 μ m wide gap between the LED and the lens. The images were 710 × 480 px² 328 in size with a resolution of $2 \mu m/px$ allowing the detection of droplets larger than $8 \mu m$ in diameter. The images were then split into two 480×480 px² square images (Figure 9 (c)); to double the number of acquired images. 329 330 Further information on the experimental setup and the fluids can be found in Khalil et al. (Khalil et al., 2010).





Figure 9. Schematic diagrams of the setup used to obtain images of a water-in-TPH emulsion. (a) Diagram of the video probe
used to capture the images of the emulsion flowing through the gap between the lens and the LED. (b) Diagram of the
double-jacketed cylindrical tank containing the emulsion, the video probe and the propeller used to stir the mixture. (c)
Examples of the raw images obtained, after binarization, and after cropping.

336 The volume fraction and the PSD were retrieved from 3 870 images. In this experiment, the volume fraction

337 could be controlled but the PSD of the dispersed phase was unknown. The predicted PSDs were therefore

338 compared with values obtained by the Hough transform (Amokrane et al., 2014). The two machine learning

and CNN^{PSD} , were retrained on a new simulated dataset corresponding to the experimental

setup, consisting of 240×240 pixel images in three channels ($240 \times 240 \times 3 \text{ px}^3$) as before.

341



Figure 10. Validation results for a liquid–liquid system (water-in-TPH emulsion). (a) Examples of the predicted volume
 fraction for different experimental volume fractions (the droplet density in the images), (b) Comparison of predicted and
 Hough transform PSDs (PSD vs PSD_{HT}).

Figure 10 (a) shows the predicted hold-up values obtained beneath corresponding experimental images. The values are consistent with the actual number of droplets over a wide range of hold-up values. The mean predicted hold-up of 4.74% is in agreement with the expected value of 5% (RE = 0.052). Figure 10 (b) compares the mean predicted PSD computed from the 3 870 images to the distribution obtained using the Hough transform. The two distributions match closely ($\overline{MAE} = 0.012 \pm 0.047$), but the CNN predicts greater proportions of larger particles, suggesting that it performs better than the HT for large particles and particles that intersect with the image boundaries.

353

5. Conclusion and perspectives

This article describes a new image processing technique for dense flow imaging, combining a stochastic geometrical model with convolutional neural networks to retrieve 3D properties of particle systems using only 2D projected images. Using this approach, the volume fraction of the dispersed phase and the particle size distribution are predicted directly from the images. A stochastic model was developed to generate training 359 images labeled with the targeted 3D properties of the system. Two CNN architectures were built, one to quantify 360 the volume fraction of the observed particle field and the other to recover the PSD. After numerical validation on 361 simulated images, the approach was validated on experimental images of calibrated beads and of an emulsion in 362 a tank. Results confirm that average 3D quantities can be measured using this approach from a single optical 363 access, even for flows with high-dispersed phase volume fractions. In particular, this new approach remains 364 effective in the presence of high particle overlap, where traditional image-processing techniques typically fail. 365 The approach is also computationally efficient, with the analysis of thousands of images typically completed in 366 less than 30 min. Moreover, the method is easy to implement and can readily be re-purposed by transfer learning 367 to study different types of multiphase flows, making it suitable for many applications in chemical engineering. 368 These encouraging results highlight the value of combining a stochastic geometry model with deep learning to 369 predict 3D information. Note that while a lognormal function was used for the PSD in this study, the model can 370 be trained to predict any type of discrete distribution. The model could also be extended to study ellipsoidal 371 particles, which are more representative of gas-liquid flow. This should be relatively straightforward since 372 Matern type II models for ellipsoidal shapes have already been described in the literature (de Langlard et al., 373 2018b).

Future work will focus in the short term on improving the stochastic model by including spatial heterogeneity
and more complex shapes (e.g. cap-shaped or skirted particles); to better account for spatial inhomogeneities in
the flow images and the variety of particle shapes typically encountered in dispersed phases of multiphase flows.
Also, more complex CNN architectures will be studied in order to improve the results and treat irregular shapes
(Ma et al., 2022; Murata et al., 2020).

379 List of symbols and abbreviations

Symbols

W	Observation window (m^2)
l	Length of the projection direction (m)
r_i	Radius of the <i>i</i> th particle
t_i	Time of birth of the <i>i</i> th particle
n_p	Total number of particles
n	Number of outputs for CNN ^{PSD}

Number of training images
Number of validation images
Number of testing images
Total number of generated images $(N = N_t + N_v + N_e)$
Truncated Lognormal law
First support of the truncated law
Second support of the truncated law
Number of particles in the <i>j</i> th class
Diameter of the <i>j</i> th class (m)
Total volume of the dispersed phase in the <i>i</i> th image (m^3)
Total volume of the mixture in the <i>i</i> th image (m^3)
Volume fraction of the <i>i</i> th image
Predicted volume fraction of the <i>i</i> th image
Particle size distribution of the <i>i</i> th image
Predicted particle size distribution of the <i>i</i> th image
Variance
Mean diameter (<i>m</i>)
Synthetic particle image
Binarized particle image
Surface fraction
Overall Mean Absolute Error

Greek Symbols

γ	Surface tension (N/m)
η	Viscosity (Pa.s)
λ	Point process intensity
μ, σ	Lognormal parameters
ρ	Density (kg/m ³)
ϕ	Hold-up or Volume fraction of the dispersed phase

Abbreviations

3D	Three-dimensional
2D	Two-dimensional
CNN	Convolutional Neural Network
CNN [¢]	Network for predicting volume fraction
CNN ^{PSD}	Network for predicting particle size distribution
DL	Deep Learning
DIH	Digital In-line Holography
FC layer	Fully Connected layer
FIB-SEM	Focused Ion Beam-Scanning Electron Microscope
FBRM	Focused Beam Reflectance Measurement
MAE	Mean Absolute Error
OMAE	Global Mean Absolute Error
PMMA	Poly (Methyl Methacrylate)
PSD	Particle Size Distribution
RE	Relative Error
std	Standard Deviation
IoU	Intersection over Union

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